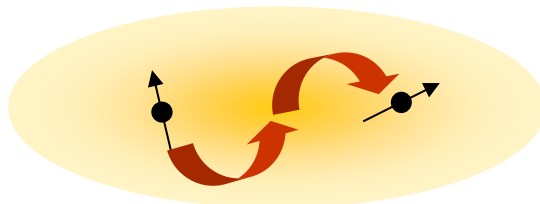


NMR spectroscopy

- J – interaction / J – coupling
- Splitting patterns
- Homonuclear and heteronuclear couplings
- Typical values, Karplus equation

J interaction

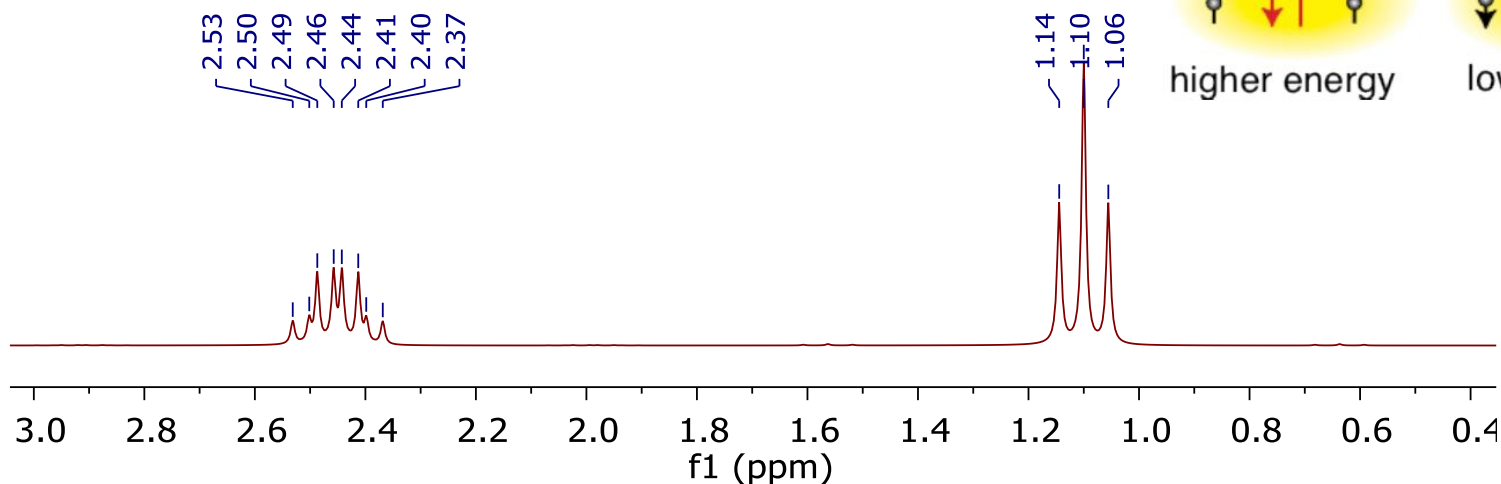
Indirect spin-spin interaction



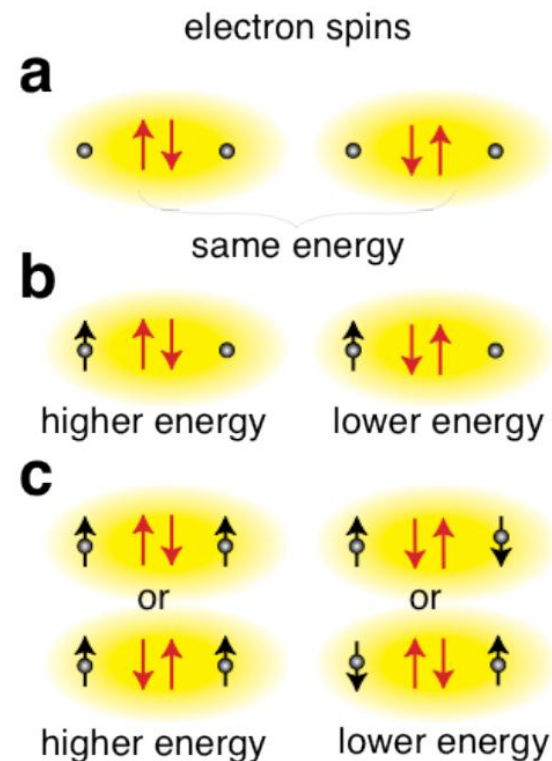
interaction mediated by common Electrons

nucleus „feels “ spin state of nearby nuclei

- Causes fine splitting of signals
- Carries information about nearby NMR-active nuclei
- Magnitude ***DOES NOT DEPEND*** on magnetic field



scalar coupling



Quantum view

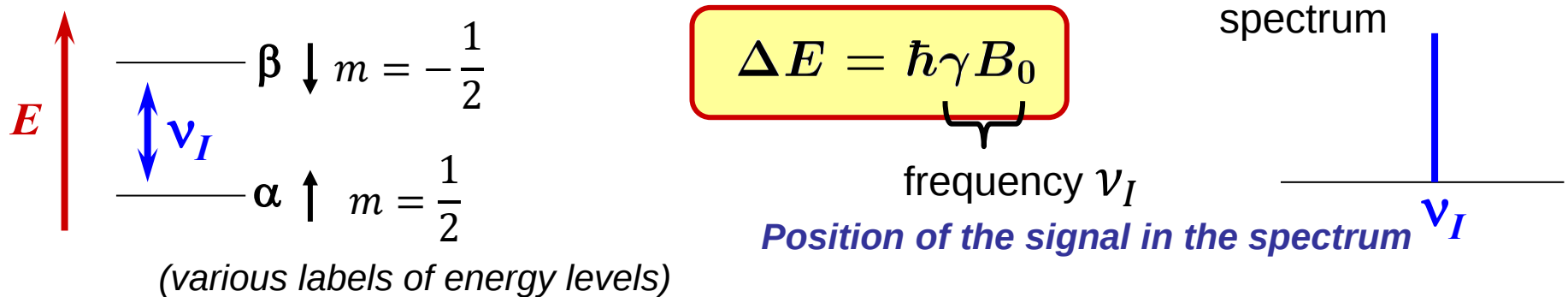
Energy of magnetic moment in magnetic field

$$E = -\vec{\mu} \cdot \vec{B} = -\mu_z B_0 = -\gamma I_z B_0 = -m\gamma\hbar B_0$$

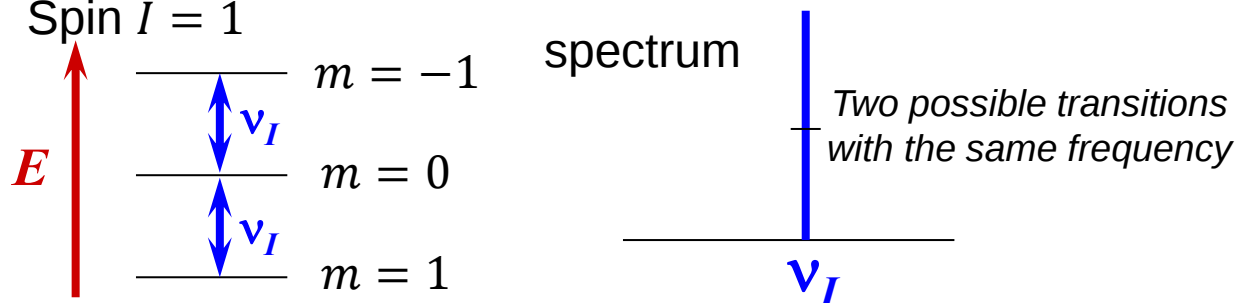
Projection of nuclear angular momentum I_z is quantized: $2I + 1$ values

$$m = (-I), (-I + 1), \dots (I + 1), I$$

Spin $I = 1/2$



Spin $I = 1$



Spin $I = 3/2$ has levels

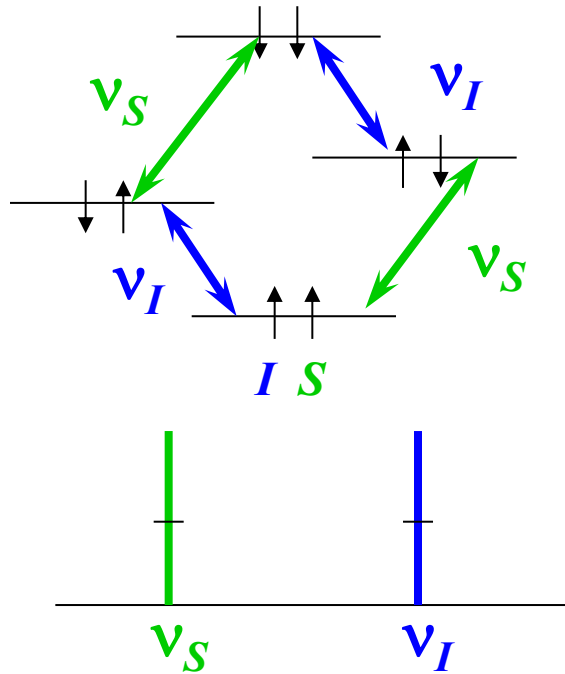
$$-3/2, -1/2, +1/2, +3/2$$

J interaction

Two nuclei with spin 1/2

Magnitude of interaction: J [Hz]

No interaction



both transitions are equal

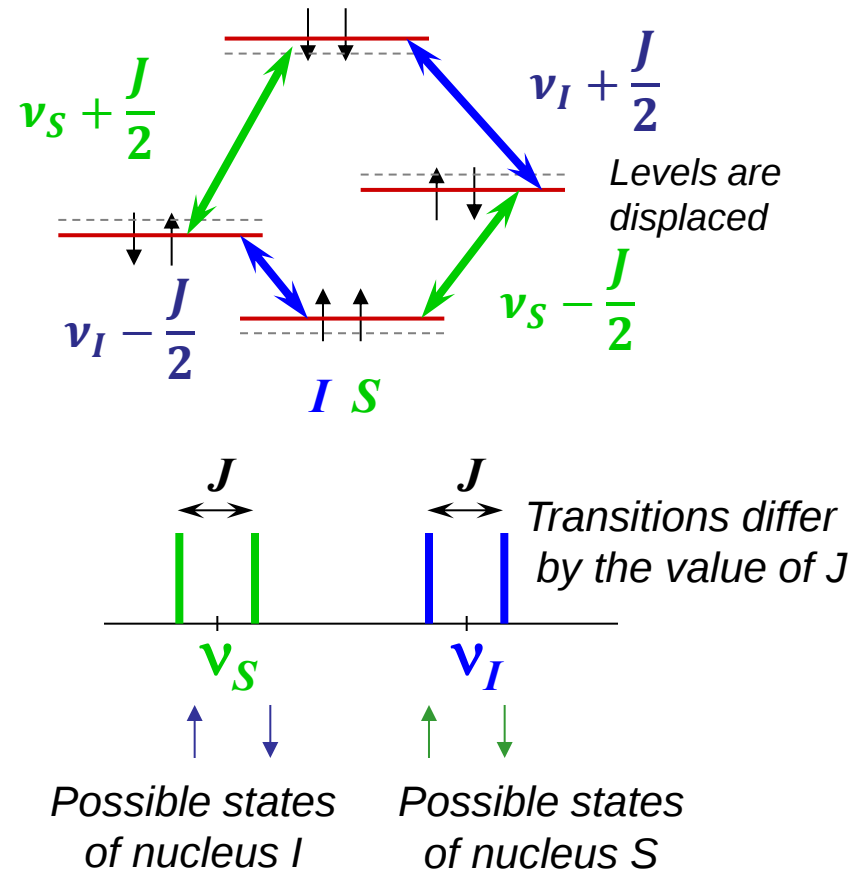
No matter of the spin state of the second nucleus

Energy contribution

$$H = 2\pi J \vec{I} \cdot \vec{S}$$

$$E_J = h J m_I m_S \quad m_I, m_S = \pm \frac{1}{2}$$

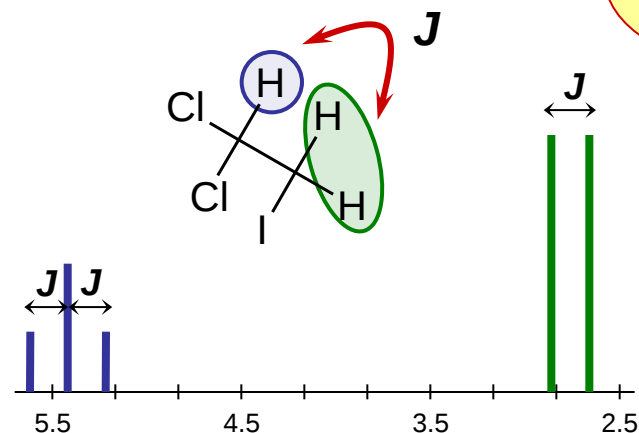
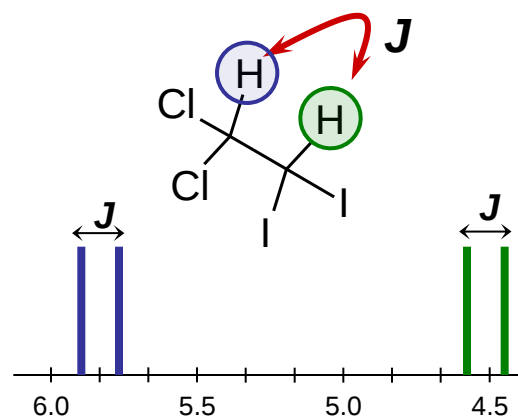
With J-couplings



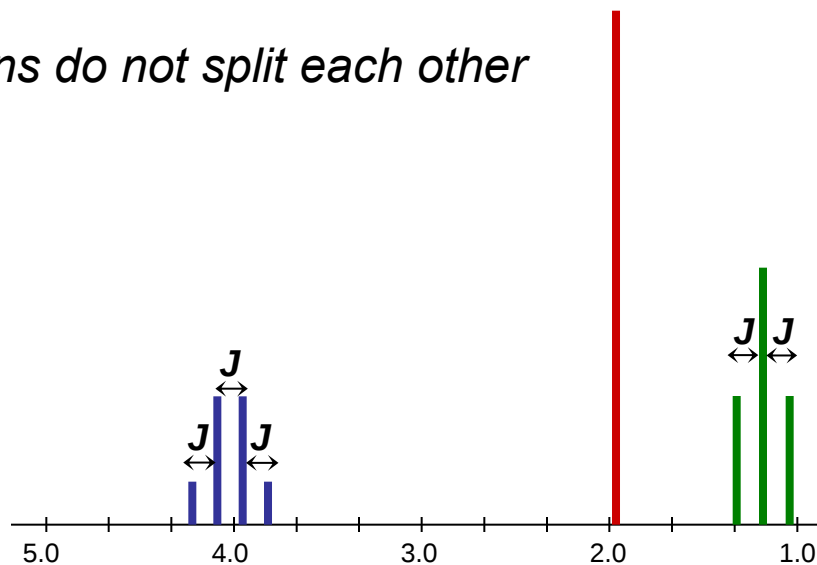
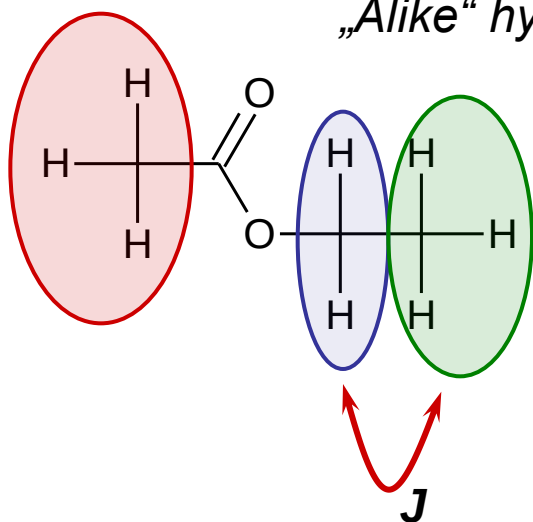
J-coupling in ^1H spectra

Signals are split by hydrogens that are „different“

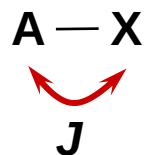
Restricting
to $^3J_{\text{HH}}$



„Alike“ hydrogens do not split each other

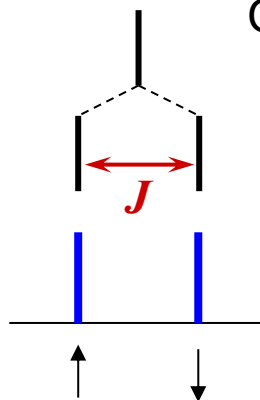


J-coupling and system of splitting



Graphical representation
of splitting

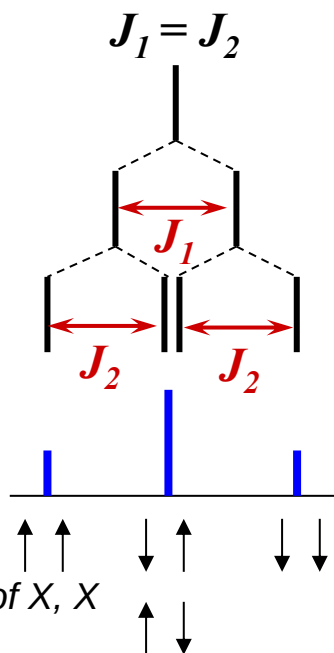
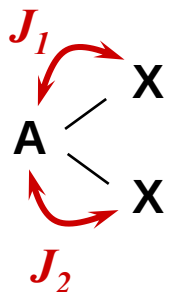
Assuming all nuclei
have $I=1/2$



spectrum of nucleus A is a **doublet**

d

Spin-states of X



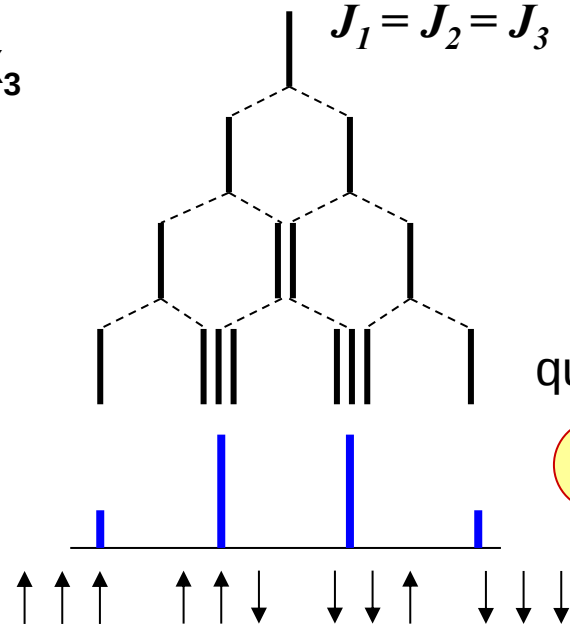
t

triplet

Spin-states of X, X



$$J_1 = J_2 = J_3$$



quartet

q

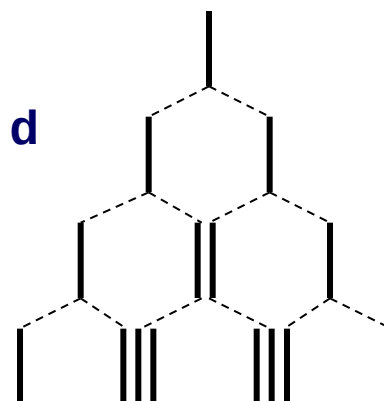
Spin-states of
X, X, X

J-coupling and system of splitting

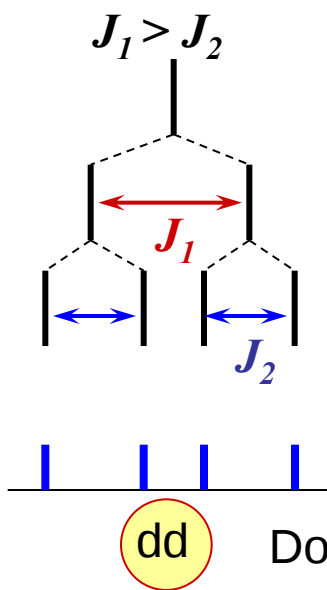
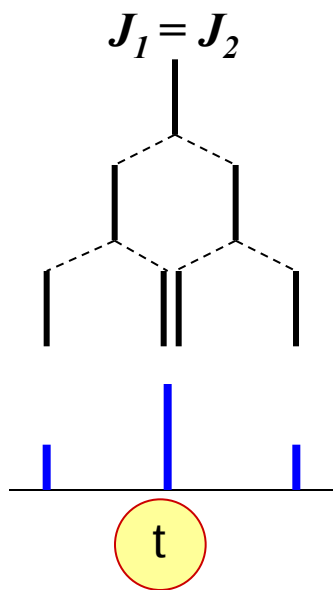
Splitting by n equivalent Hydrogens

n	1				singlet s			
1	1	1			Doublet			
2	1	2	1			triplet t		
3	1	3	3	1			quartet q	
4	1	4	6	4	1			pentet p

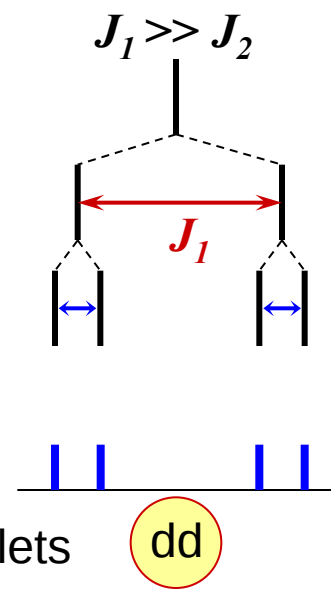
Pascal's triangle



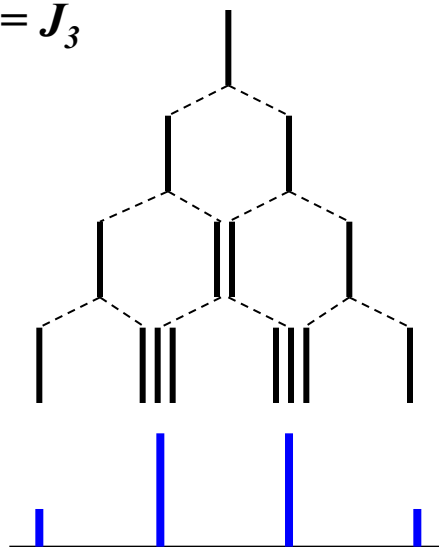
*Rule of
 $n+1$ lines*



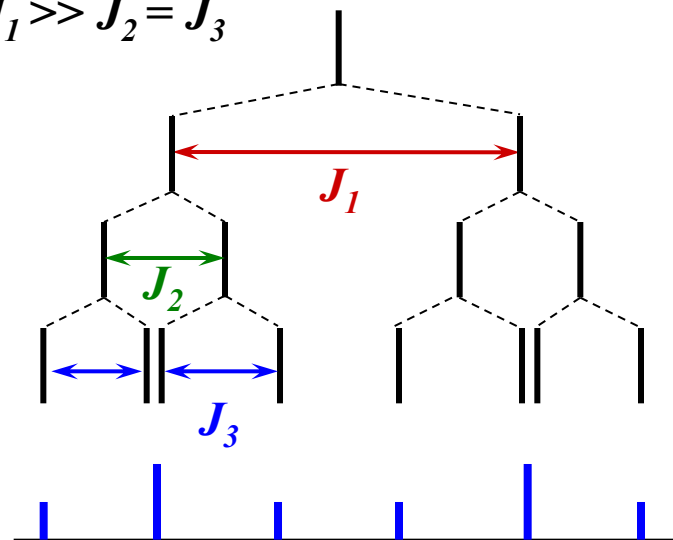
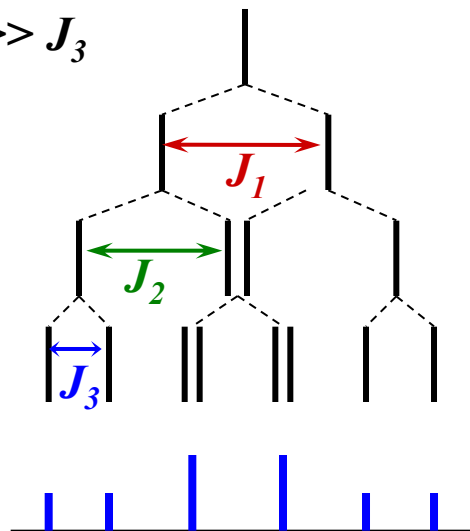
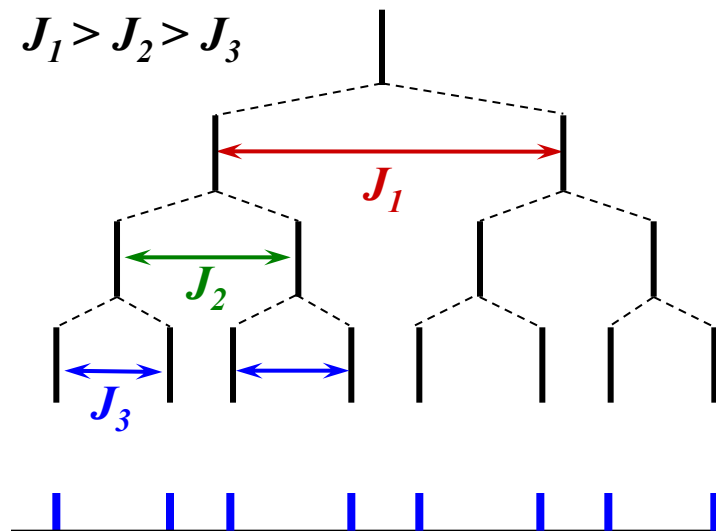
Doublet of doublets



J-couplings – structure of multiplets

$$J_1 = J_2 = J_3$$


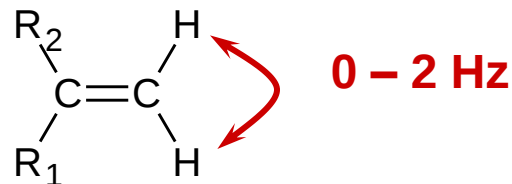
q

$$J_1 \gg J_2 = J_3$$
 dt
$$J_1 = J_2 \gg J_3$$
 td |
$$J_1 > J_2 > J_3$$


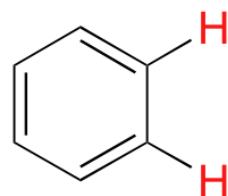
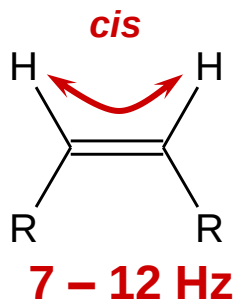
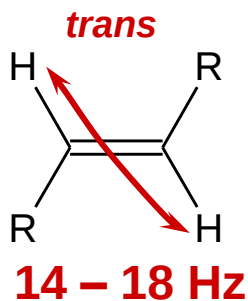
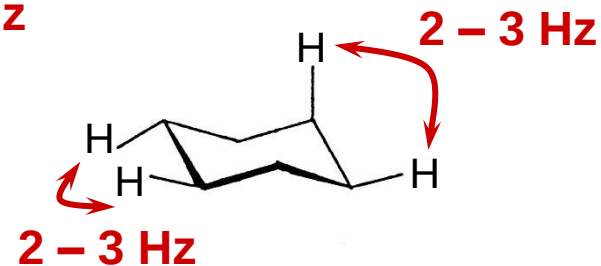
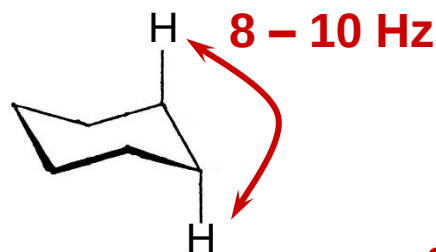
ddd

J-couplings in hydrogen spin systems

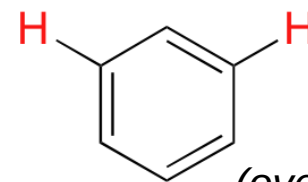
Geminal $^2J_{HH}$ (Through two bonds)



Vicinal $^3J_{HH}$ (Through three bonds)



ortho-benzylic:
 $^3J = 6-10 \text{ Hz}$

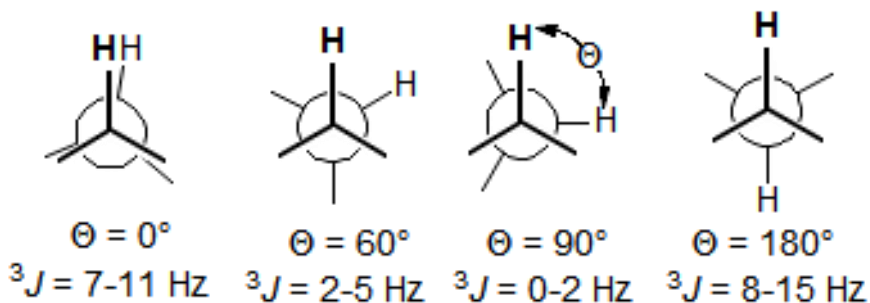
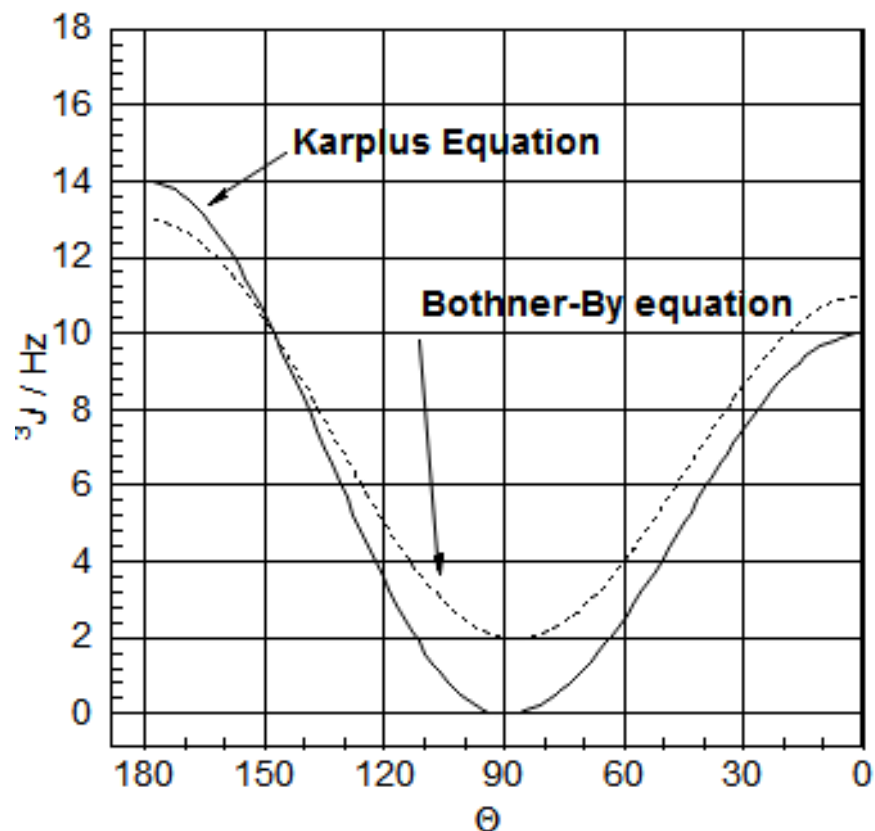


meta-benzylic:
 $^4J = 0-4 \text{ Hz}$

(over 4 bonds)

Karplus equation

Dependency of $^3J_{HH}$ on the dihedral angle



Karplus Equation

$$^3J_{HH} = J_0 \cdot \cos^2 \Theta - K$$

$$J_0 = 14 \text{ (90-180°), } J_0 = 10 \text{ (0-90°), } K = 0$$

Bothner-By equation

$$^3J_{HH} = 7 - \cos \Theta + 5 \cdot \cos 2\Theta$$

This empirical dependence of the size of the J-coupling can be used to determine the conformation of the molecule

Heteronuclear J-interaction

$^{13}\text{C}-\text{H}$
Spin-states of H

$m = -1/2$ ↓
 $m = +1/2$ ↑

Doublet
1 : 1

$^{13}\text{C}-\text{H}_2$

↓ ↓
↓ ↑ ↑ ↓
↑ ↑

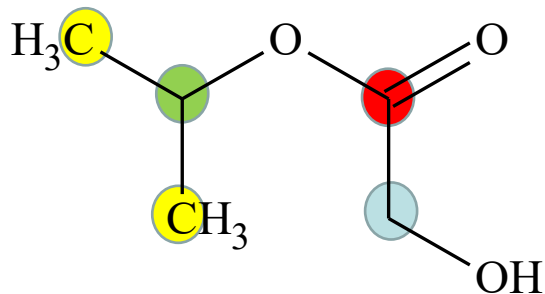
triplet
1 : 2 : 1

$^{13}\text{C}-\text{H}_3$

↓ ↓ ↓
↓ ↓ ↓ ↑ ↓ ↓ ↑ ↓ ↓
↑ ↑ ↑ ↑ ↓ ↑ ↓ ↓ ↓
↑ ↑ ↑

quartet
1 : 3 : 3 : 1

^{13}C spectrum

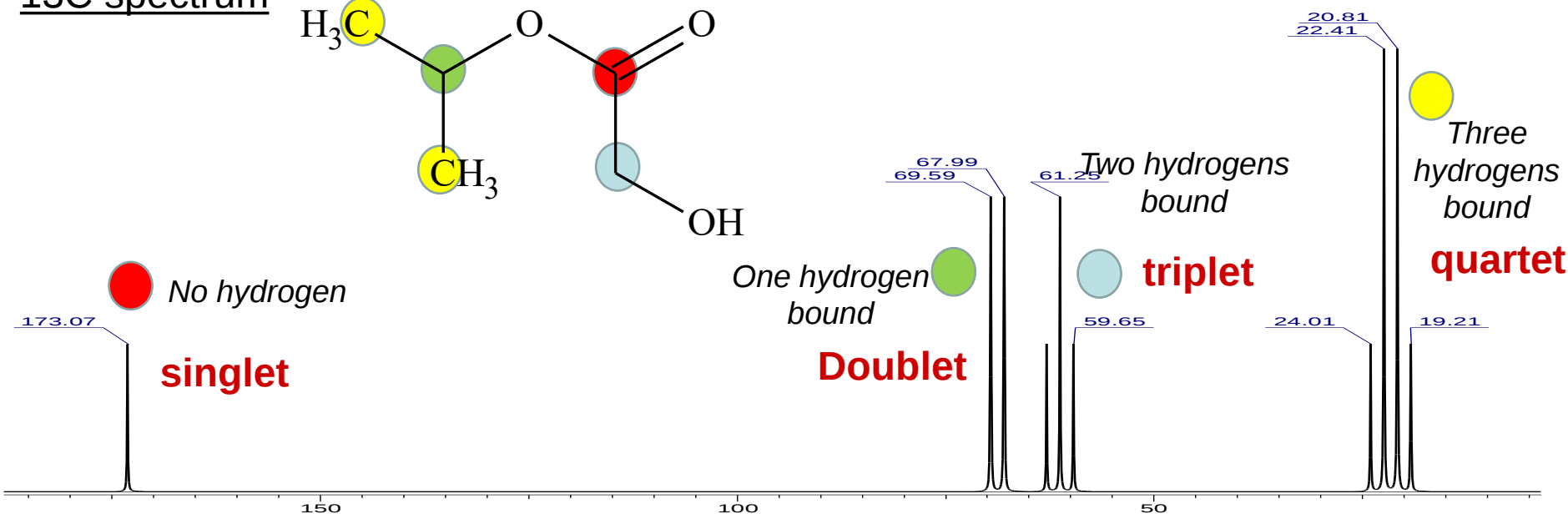


173.07
singlet
No hydrogen

One hydrogen bound
Doublet

Two hydrogens bound
triplet

Three hydrogens bound
quartet

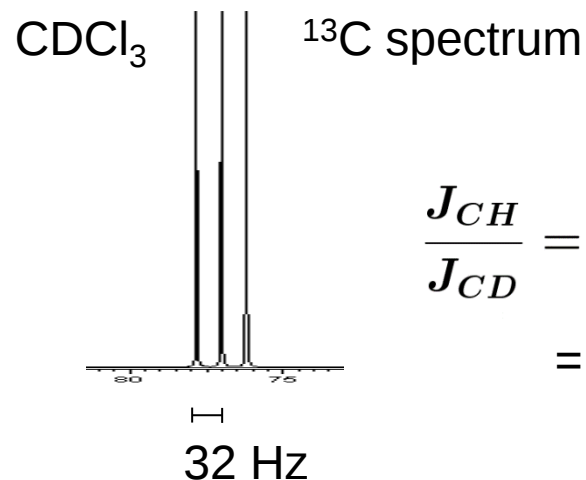
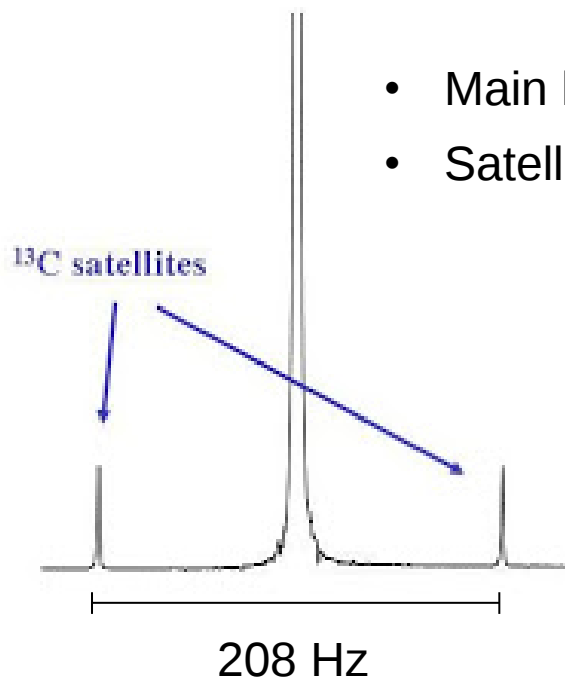


Heteronuclear J-interaction

	H ₃ C-CH ₃	H ₂ C=CH ₂	C ₆ H ₆	HC≡CH
$^1J_{CH}$	124,9	156,4	158,4	249,0

$^2J_{CH}$ and $^3J_{CH}$ is about 8 Hz

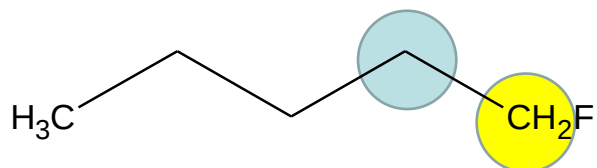
1H spectrum CHCl₃



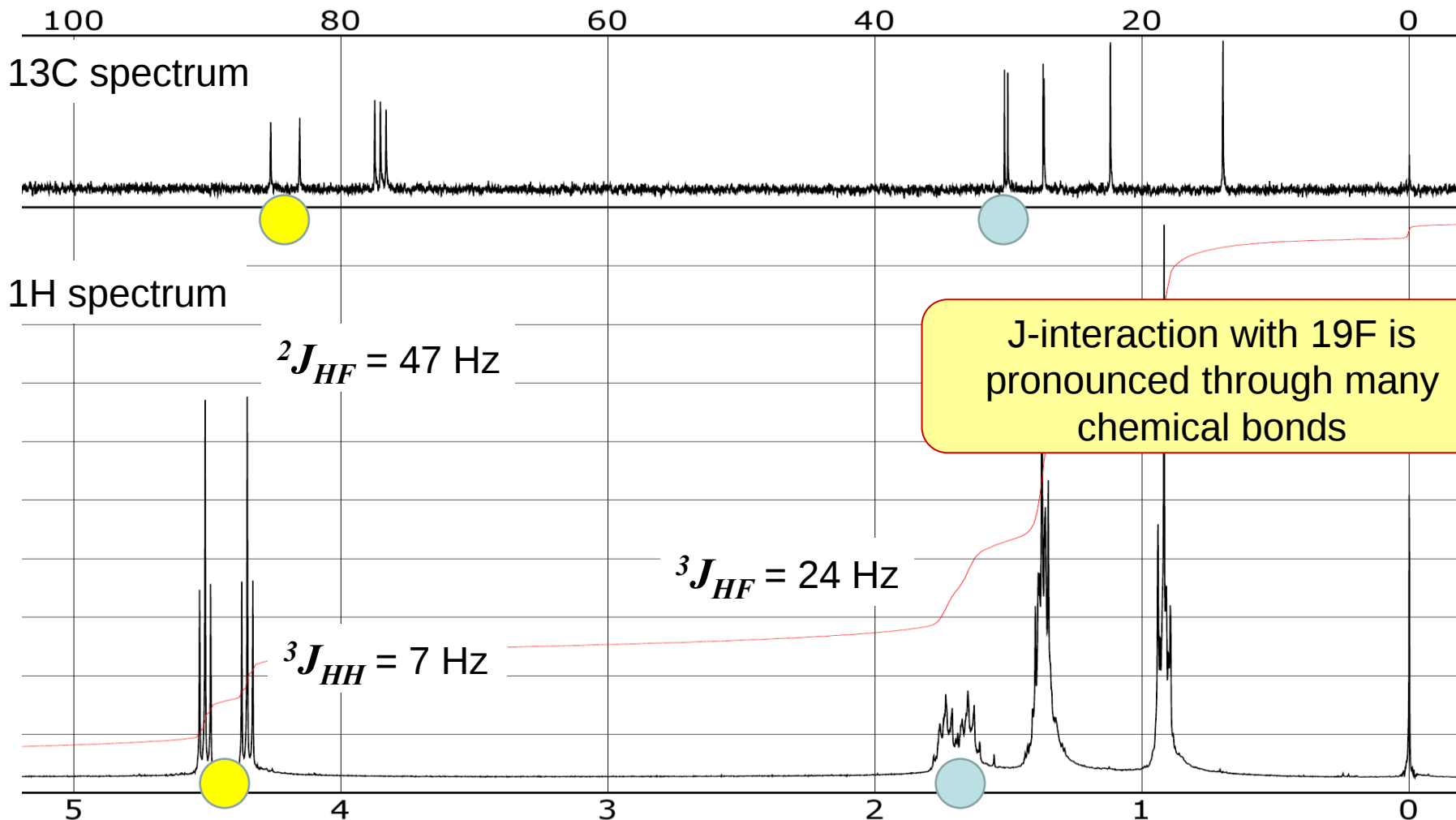
$$\frac{J_{CH}}{J_{CD}} = \frac{\gamma_H}{\gamma_D} = 6.51$$

J interaction with other nuclei: ^{19}F

$$I = 1/2$$

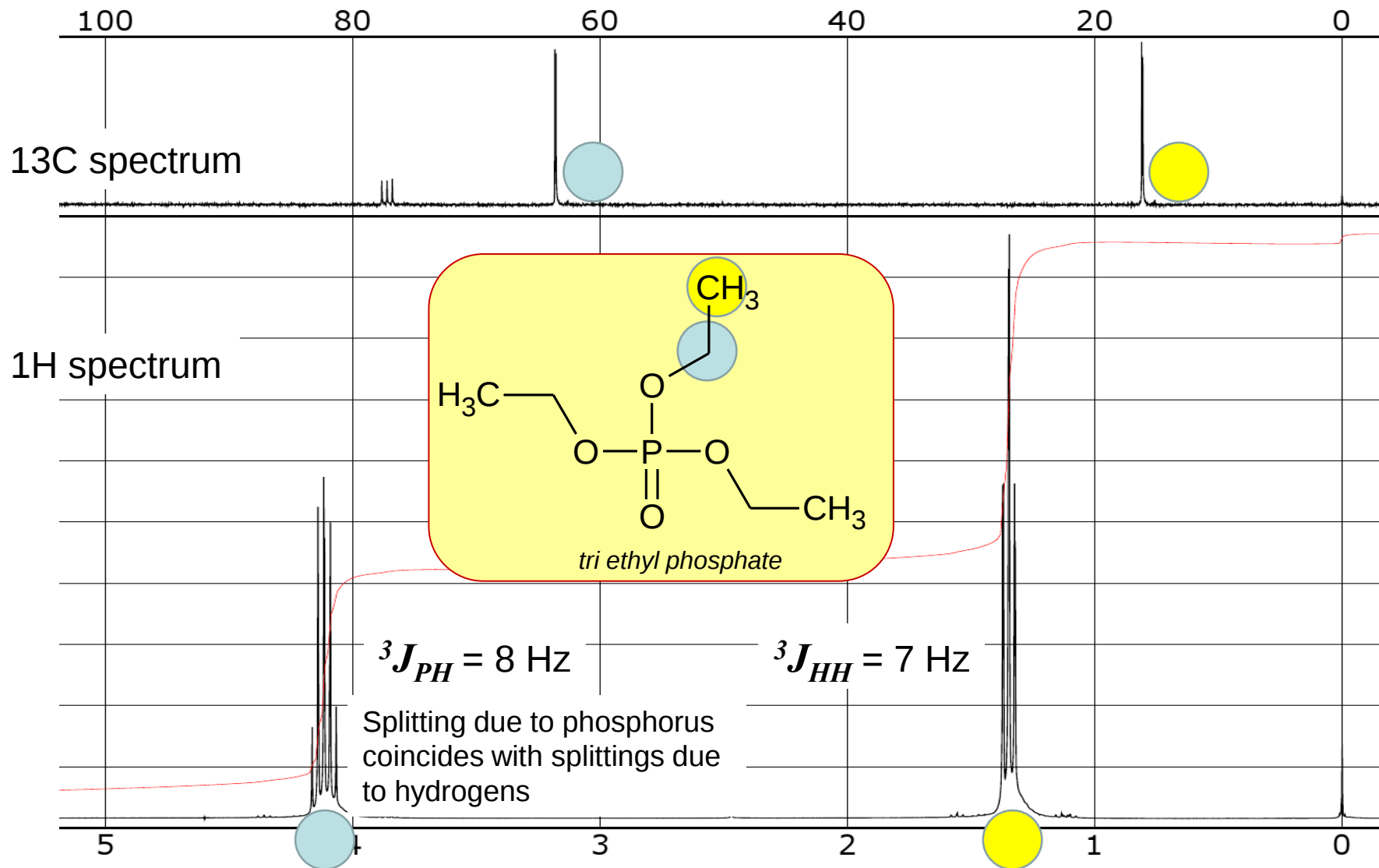


$$^1J_{\text{CF}} = 167 \text{ Hz} \quad ^2J_{\text{CF}} = 26 \text{ Hz} \quad ^3J_{\text{CF}} = 7 \text{ Hz}$$



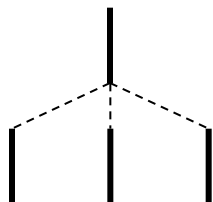
J interaction with other nuclei : ^1H - ^{31}P

$$I = 1/2$$



J interaction with other spins

AX

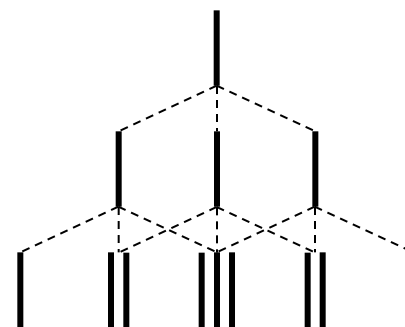


1 : 1 : 1

CDCl_3

... applies to hydrogens in NH_4^+ too

AX₂



1 : 2 : 3 : 2 : 1

CD_2Cl_2

When X has spin $I = 1$

possible states

-1, 0, +1

^{11}B (80%) : spin-3/2

Possible states

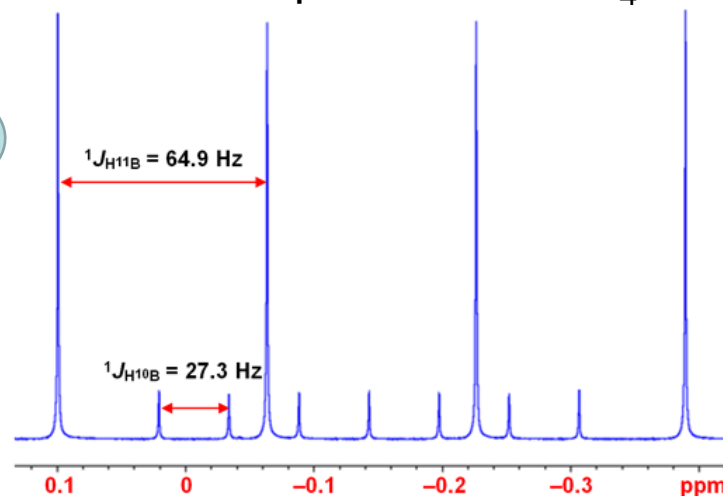
-3/2, -1/2, +1/2, +3/2

Boron has two
naturally occurring
NMR active isotopes

^{10}B (20%) : spin-3

Possible states

Proton spectrum of $^+\text{BH}_4$



Magnetic equivalence

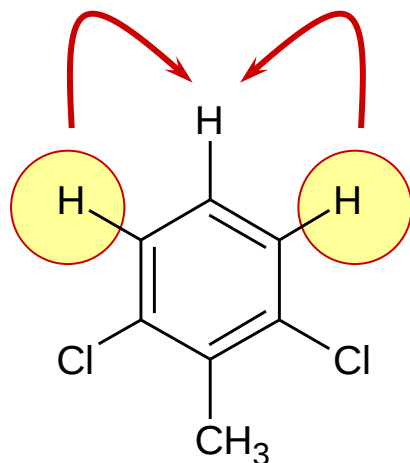
Nuclei are magnetically equivalent if:

- are chemically equivalent

and at the same time

- have the same geometric configuration in relation to all other NMR active nuclei in the molecule (that is, have the same J-interaction)

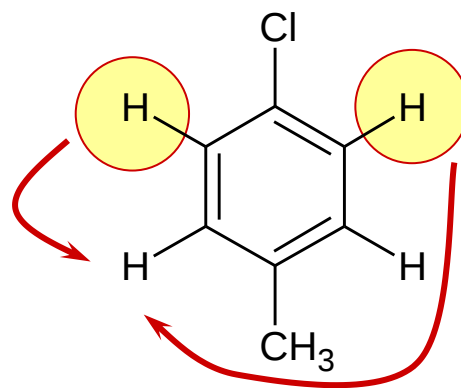
Equal configuration



Magnetically equivalent

J-interaction between these protons is not manifested in the spectrum

Different configurations

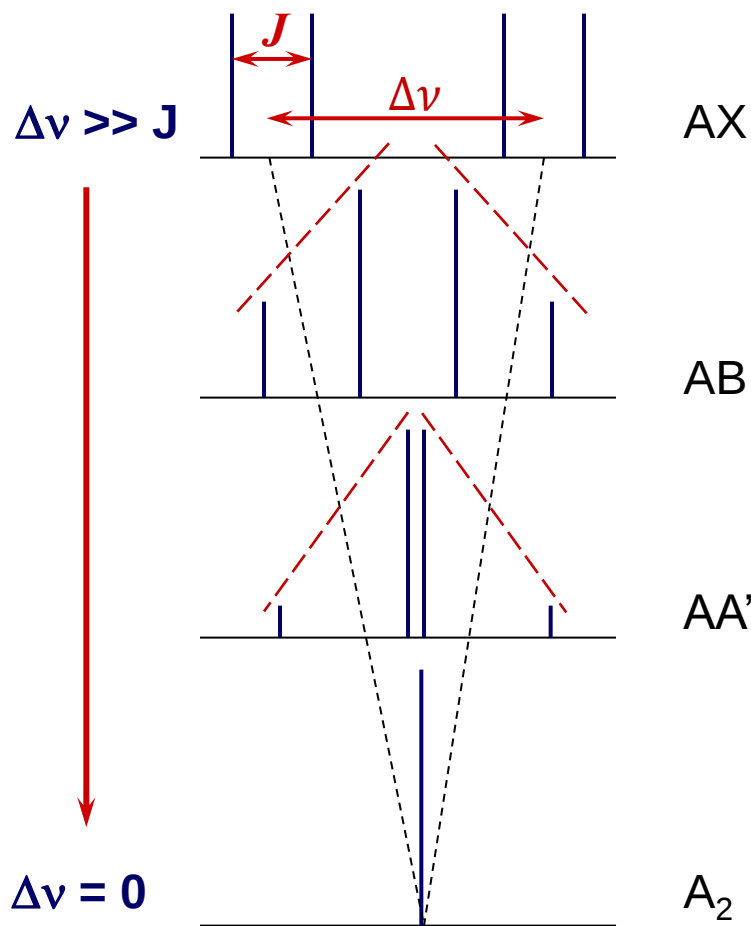


**Chemically equivalent,
but magnetically non-equivalent**

Their J-interaction causes complicated splitting pattern

Order of the spectrum

Depends on magnitude of J-coupling and difference of chemical shifts $\Delta\nu$
(both in units of Hz)



1st order spectrum
(1st order perturbation theory)

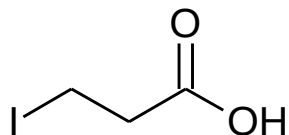
"Pseudo-First Order"

higher-order spectrum

Roof effect

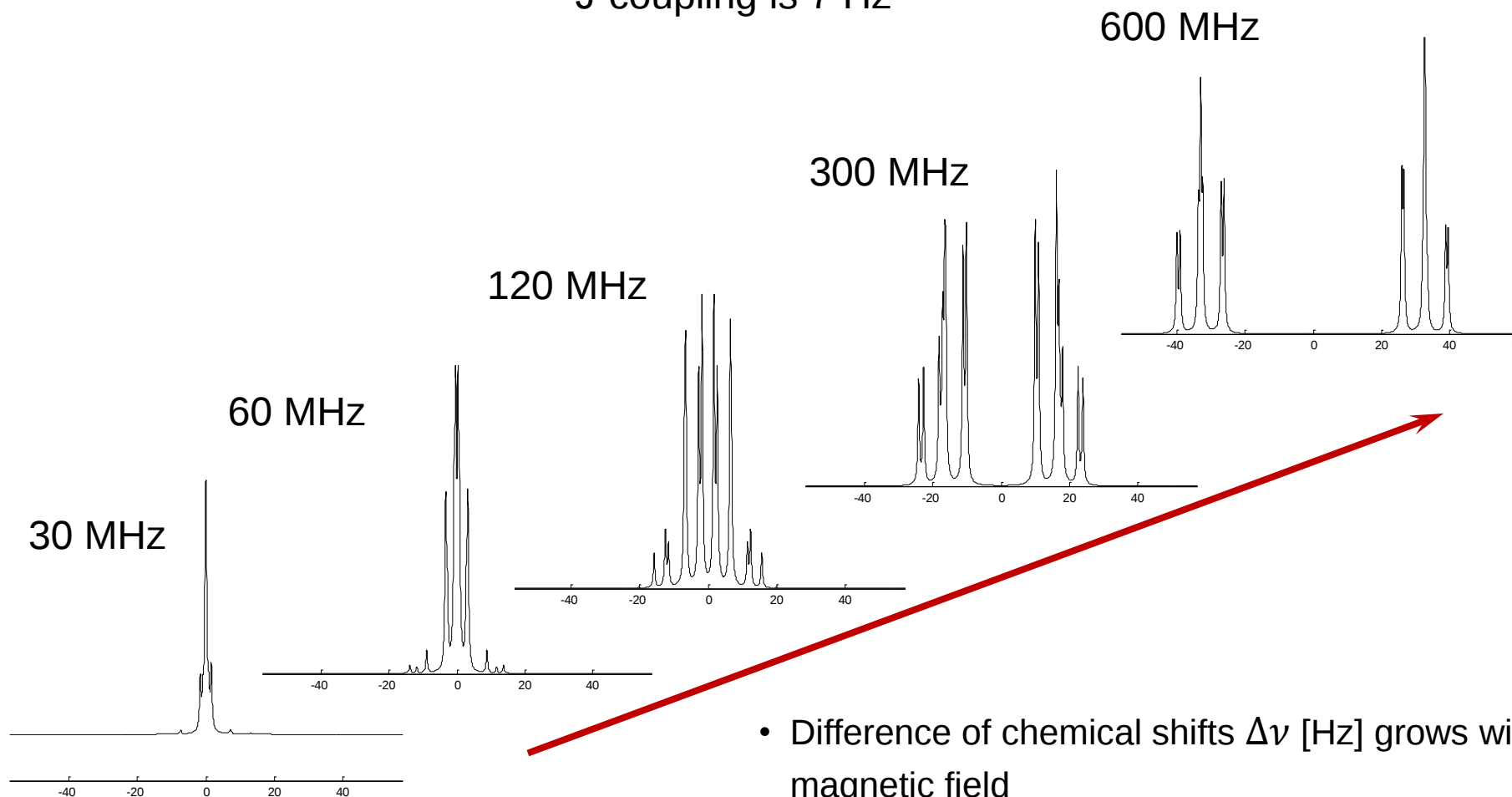
- The intensities of the lines in the doublet do not correspond to the ratio 1:1
- Inner lines (towards the coupling partner) have a higher intensity
- This can be used in the interpretation of spectra

Order of the spectrum



Chemical shift difference of 0.11 ppm

J-coupling is 7 Hz



- Difference of chemical shifts $\Delta\nu$ [Hz] grows with magnetic field
- J-coupling is independent of the magnetic field